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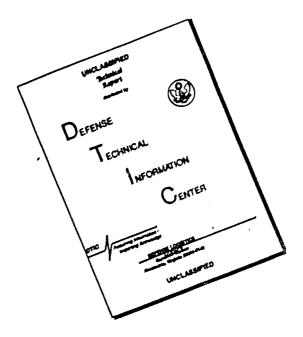
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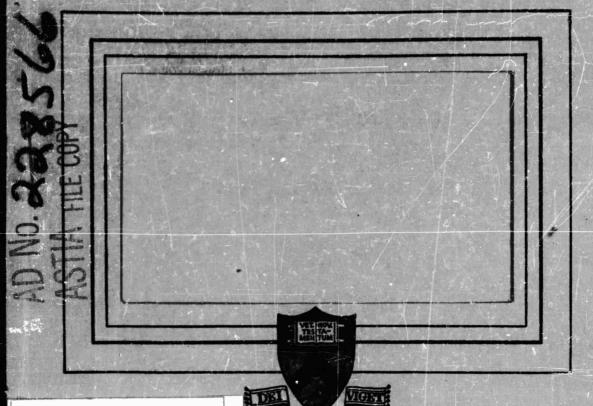
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PRINCETON UNIVERSITY

DEPARTMENT OF AERONAUTICAL ENGINEERING

Addenda and Errata to

METAL COMBUSTION PROCESSES

by Irvin Glassman

ARS Preprint No. 938-59

AFOSR T.N. No. 59-1043

Princeton University Aero Eng'g Lab Report No. 473

- Pg 6, line 17 After the word "oxidizer" insert "at the flame front".
 - line 18 After the word "oxidizer" insert "at the front".
 - line 24 After the word "flames" insert the following sentence.

 "The fact that the oxide condenses simplifies the diffusion equations generally used in droplet analyses."
 - line 24 Insert the word "the" before "metal".
 - line 26 Change last word to "Thus".
- Pg 7, line 2 Insert word "curve" after "distribution".
 - line 19 After "metals" insert "or metallized solid propellant".
- Pg 8, line 19 After the number "1100°K" insert the following sentence in parenthesis "(At this temperature the vapor specie above a B-B₂0₃ mixture is more correctly written as the complex molecule B₂0₂ (93))".
- Pg 9, line 14ff Change stoichiometry of the equation to

4/3 Al(1) + 1/3 Al2 O3(5) -> Al2O(g)

The paragraph following should read: "Farber lists the value of $4\,{\rm H_f}^{\rm O}$ (25°C) for Al₂O (g) as -34 kcal/mole. Since the $4\,{\rm H_f}^{\rm O}$ of Al₂O₃ (S) is -399 kcal/mole, the above reaction is highly endothermic and thus limits the temperature. Since BO has a $4\,{\rm H_f}^{\rm O}$ of only -5.3 kcal/mole a similar situation exists for boron. It is obvious for metal oxides which decompose into elemental metal vapor, metal suboxides and/or oxygen atoms."

- Pg 10, line 8 After end of sentence insert the following sentence.

 "A list of the major species for the vaporization of metal oxides under various conditions is given at the end of Ref. 93."
 - line 17 Insert "adiabatic" before "temperature".
 - line 18 Insert "in oxygen" after "aluminum".
- Pg 12, line 10 Insert "and physics" after "chemistry".
- Pg 14, line 10 Insert "burning" before "droplet".
 - line 15 Last word should read "partial".
 - line 23 After "coefficient" add "relative to the metal".
- Pg 15, line II Should read ".....reduces to a more simplified equation....." instead of ".....to the equation.....".
- Pg I7, line I4 Should read ".....pressure of the oxygen present, that is $P_{0_2,B}$ " instead of ".....that is $P_{0_2,B}$ "
 - line 17 Change "of" to "on".
 - line 20 Change "good" to "valid".
- Pg 18, line 17 Insert "generally" before "is".
 - line 19 Should read "......should have flame fronts close to....." instead of ".....should be close to.....".
 - Insert the following sentence. "In fact, if one assumes that the thermal conductivity of the metal vapors are approximately the same, as a rough approximation, one could postulate that in the initial burning phases that the separation distances between the flame front and the metal surface for aluminum and sodium flames are relatively of the same order; that beryllium flames have a separation distance a factor of 4 less than that of aluminum; lithium a factor of 2 less; magnesium a factor of 3 greater."

Pg i9, Table III Heading of Column 7 should read "TM". Add a column 12, with heading and values as follows:

· ·	V V T _{MO} -T _M
Eiement	Н
Be	0.08
LT	0.16
ΑI	0.34
Mg	i.i4
Na	0.32
₽6H6	27.0

Pg 20, Paragraph 4 At the end of paragraph 4 add the reference letter (k).

last line Add "initially" atter "but".

- Pg 23, line 18 Change "alkylated compounds" to "liquid alkylated metal compounds".
- Pg 24, line 2 Change that part of the sentence beginning "less oxide deposition....." to "less oxide deposition may or may not be a great oversimplification, for nucleation products appear to be of the order (I,m) of 0.1 to 1 \mu in diameter and the particles must diffuse a distance of the order of 100 to 1000 particle diameters."

Add the following references: "

- k. Glassman, I., "The Stability of Propellants and the Theory of Thermai Ignition", Princeton University Aero Eng'g Lab Report No. 460, "AFOSR TN 59-586, ASTIA AD No. 217 185, May, 1959.
- Setze, P. C., "Liquid Boric Oxide Particle Growth Rates in a Gas Stream from a Simulated Jet Engine Combustion Chamber", NACA EM 55I 20a, April 30, 1957.
- m. Schreier, S., "The Effect of a Condensed Exhaust Phase on Rocket Performance", M.S.E. Thesis in Aero Eng¹g, Princeton University, 1956.

Add the following to the bibliography:

93. Inghram, M. G., and Drowart, J., "Mass Spectroscopy Applied to High Temperature Chemistry", Advance Papers of an Internationa! Symposium on High Temperature Technology, Asilomar, Calif., Oct. 6-9, 1959, arranged by Stanford Research Institute, Menlo Park, Calif. This manuscript contains 137 references, many of which are pertinent to metal combustion phenomena.

- 94. Kingery, W. D., "Oxides for High Temperature Applications", ibid.
- 95. Ackermann, R. J., and Thorn, R. J., "Reactions Yielding Volatile Oxides at High Temperatures", XVI Congrès International de Chemie Pure et Appliquée, Paris, 1957.

United States Air Force Air Research and Development Command Air Force Office of Scientific Research

A Technical Report on Contract AF 18(600)-1527

METAL COMBUSTION PROCESSES

by Irvin Glassman 377

Aeronautical Engineering Laboratory Report No. 473

August 1959

Department of Aeronautical Engineering
PRINCETON UNIVERSITY
Princeton, New Jersey

AP STRACT

Since the light metallic elements undoubtedly will play a greater and greater role as propellants in advanced chemical and nuclear-chemical propulsion systems, interest in their burning characteristic is increasing.

Presented here are some preliminary conclusions on burning characteristics.

These conclusions are based on fundamental physical considerations and not on experimental results. An analytical approach to calculate the burning rate of metals is also suggested. This approach differs from the diffusion-droplet approach in that it includes radiation feed-back and loss terms. Such terms can be important at the high temperatures of the diffusion film surrounding a burning metal.

The verification of many of the postulates given in this report can be carried out ideally at some later date on the high pressure double rocket motor flow reactor developed under the subject contract.

METAL COMBUSTION PROCESSES

Irvin Glassman² Princeton University, Princeton, New Jersey

I. INTRODUCTION

Interest in artificial light sources for photography and metallic dust expiosions stimulated early work on the combustion of metals. Recent endeavors to obtain higher performance with rocket motors (a)3 and ram engines (b) have led many propellant engineers to formulations which contain certain metallic elements in some form. The addition of aluminum powders to solid propellant combinations has given high specific impulse and better operating characteristics (a). However, aluminum is not necessarily the only metallic element of interest to the propulsion engineer. An interesting presentation of which substances may be best as propellants has been given by Carpenter (c) and is shown in Figure i. In this figure there is piotted the heat of combustion in Btu per pound of fuel and oxidizer versus the atomic number. Since the flame temperature is related to the heat of combustion and in rocket motors both fuel and oxidizer must be carried, this plot gives some idea as to the relative effectiveness as propellants of various elements and compounds formed from these elements. Of course, the molecular weight of the products of combustion is important, as well, and this consideration must be taken in account. Although this graph is based upon oxygen as the oxidizer, the relative order would not be effected too greatly by selection of another oxidizer. One should note that most compounds will give a heat of combustion

i. The author's work in combustion is supported by the United States Air Force Office of Scientific Research of the Air Research and Development Command under Contract AF 18(600)-i527.

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^{3.} Lower case letters appearing in parentheses refer to articles listed in the references at the end of this manuscript; numbers appearing in parenthesis refer to articles appearing in the bibliography which follows the references.

between the points representing the elements from which they are constituted. Immediately, one can see for propeliant considerations what compounds containing elements with atomic numbers greater than 14 are of no concern, for such compounds could not compete with gasoline, and their products of combustion undoubtedly would have high molecular weights. The prominence of certain of the metallic elements as propellants is accentuated by a plot of this type and the fact that many have a higher heat of combustion than hydrogen is not generally realized. It appears that the use of the elements, or compounds containing lithium, boron, berryllum, aluminum, magnesium, and sill-con as propellant additives could be most attractive. However, although these propellants will give high flame temperatures, their burning rate and other physical characteristics must be such that the efficiency and operating characteristics of the propulsive device are not impaired.

Figure 2, also taken from Carpenter (c) makes a similar comparison for air breathing engines in which the heat of combustion in Btu per pound of fuel is most important. From this figure likely fuels could be formulated from slurries of the elements, or compounds containing boron, aluminum, and magnesium. Much of the same considerations discussed above hold here as well.

A relative rating of various compounds based on the heat of combustion per pound of fuel and per pound of fuel and oxidizer is given in Table I. Hydrogen which has the highest heat of combustion per pound of fuel ranks sixth when compared on the basis of Btu per pound of fuel and oxidizer. However, the molecular weight of the products of hydrogen rocket combinations is low compared to the effective molecular weight of the products of the other substances and thus hydrogen fares much better on a specific impulse comparison. Actually Table I simply specifies which metals would be most likely as additives to hydrogen and other CHN fuels. In conventional systems metals themselves as fuels would give prohibitively high molecular weight

TABLE I
HEATS OF COMBUSTION OF VARIOUS SUBSTANCES IN OXYGEN

	Fuel	Product	Btu/ b fuel	Fuel	Btu/Ib (<u>fuel + oxygen</u>)
1.	H ₂	н ₂ о (Д)	61,000	Be	10,550
2.	Be	Be0	29,400	L1	8,640
3.	В	B ₂ 0 ₃	25,400	В	7,900
4.	CH ₄	со ₂ , н ₂ о (Д)	24,000	Al	7,080
5,	C ₂ H ₂	со ₂ , н ₂ о (Д)	21,400	Mg	7 ,00 0
6.	LI	L120	18,500	H ₂	6,780
7.	С	∩ 0 2	14,100	Si	6,300
8.	SI	s10 ₂	13,500	C2H2	5 ,250
9.	Al	A1203	13,400	CH ₄	4,800
10.	Mg	MgO	11,600	Ti	4,550
11.	Ti	T1203	6,820	C	3,840

performance can be obtained with the addition of the metallic elements and thus increased attention should be given to the combustion of metals and related processes.

Undoubtedly metals added to solid propellants are consumed by diffusion flames. It can be postulated, and more will be said about this subject later, that the flames surrounding the burning metal particles represent temperatures much greater than the equilibrium adiabatic combustion temperature of the complete propellant formulation. Equilibration must take place after condensation of the metallic oxide and in later flow stages of the motor chamber. Metals added by liquid slurries should follow the same pattern.

Compounds containing metallic elements may reach a more complete oxidized state by one of two procedures. The compound can be attacked and oxidized in the same way as a hydrocarbon, or the compound may decompose into the metallic element and chemical fragments. The metal particles formed then burn as diffusion flames. In all likelihood, most, but certainly not all, alkylated metal compounds or polymers will follow the latter procedure. Those compounds oxidized in the same manner as ordinary hydrocarbons are not of interest here; however, those which lead to metal diffusion flames are.

II. THE METAL DIFFUSION FLAME

The combustion of solid materials can be divided into two catagories—the metals and the non-metals. The non-metals include for the most part carbon and the solid hydrocarbons and there are no great distinguishing features that make the combustion of these materials any different from the combustion of liquids. In fact, the categories could be better defined to distinguish between those materials whose combustion products have very low

the high boiling point metallic oxide as their products and thus have flames which exhibit certain unique or distinguishing features. A liquid such as an alkylated metal compound would fall in the same category as the pure metal since this liquid would have a high boiling oxide as a product. The temperatures that can be obtained with the combustion of metals appear to be limited to the vaporization temperature of the metallic oxide product, whether the oxide decomposes on vaporization or not. This conclusion is reached since the heats of vaporization or decomposition of the oxide are equal to the heat of reaction of the metal and oxidizer at the vaporization temperature. Eventually when the products of a metal flame diffuse to a cooler zone they condense.

Of course, the vaporization and condensation temperature varies with the pressure, and for those compounds which decompose the temperature which can be reached will also be a function of the oxygen concentration. Thus, one should judge the temperature reached from the equilibrium process which represents the vaporization dissociation process, i.e.

$$M_{x} O_{y}(\ell) \Longrightarrow \chi M(j) + j O(g)$$

$$K_{p} = (M)^{\chi} (0)^{j} \left[\frac{p}{\sum_{m}} \right]^{\chi+j}$$

where M represents any metallic element, $\sum \infty$ is the total number of gaseous moles, and P is the total pressure. One sees that as the pressure rises and the oxygen (or oxygen atom since whichever exists is correct for the argument here) concentration rises the reaction must shift to the left. More liquid thus is present and therefore higher temperature can prevail before the liquid will decompose.

Diffusion flames of ordinary liquids on solids generally are described by the double film picture (d), as depicted in Figure 3. Heat diffuses from the flame front B to the particle surface A. The fuel vaporizes

or sublimes and diffuses from the surface A to the flame front B. Oxygen diffuses from some position C (which can be taken at infinity for a stagnant atmosphere), toward the front B. It is assumed that no reaction takes place until fuel and oxidizer reach B, that they meet here in stoichiometric portions and that the reaction takes place infinitely fast in an infinitely thin layer. Obviously, the infinite reaction rate concept cannot be a true one, but analyses developed from this picture have been successful in predicting many characteristics of liquid diffusion flames, including the burning rate of the fuel (e). More complex analyses (f) which include finite remaction kinetics, but ignore inflammability limit considerations, do not too greatly alter the basic picture, and in fact, predict the same burning rates as the simplified approach.

Generally there are then some distinctive differences which must be considered between a metal diffusion flame and those flames having only gaseous products. In the metal flame the temperature at the front is known and, as defined previously, taken as the vaporization temperature of the metal oxide. However, as will be shown later, the partial pressure of the oxidizer is not known. In the ordinary liquid diffusion flame, the partial pressure of the oxidizer is specified since the basic assumption is that the fuel and oxidizer meet in stolchiometric proportion and as a consequence the flame temperature at the front must be calculated.

The second distinction of the metal diffusion flame is that a condensable oxide is formed. As mentioned earlier the formation of the oxide really establishes the basic difference between metal diffusion flames considered here and other liquid or solid diffusion flames.

There is, however, another characteristic of metal flame which is worth noting and which may be very important. Generally the temperature in pure metal combustion is so high that there is ionization in the front. So there exists in the flame front a complex of ions, electrons, and neutral

species (60). This complex produces a high emissivity in the ultra violet (69). Since the peak of Planck's radiant energy distribution shifts to the ultra violet for these high temperatures, the emissive radiant energy of such a flame is very high. In most considerations of diffusion flames it is the heat which is conducted back from the flame through the fuel vapor which raises the temperature of the droplet to the vaporization temperature, vaporizes the fuel, and heats the vapor. It would appear then with high emissivities and high temperatures (particularly at high pressures which cause high oxide vaporization temperature) that radiant energy processes in metal flames could play an important role in determining the metal burning rate. The reflectivity of the molten metal surface receiver could then be of import and in fact it is possible that the surface condition of the burning metal could play an important part in determining its burning rate. There can be, of course, a radiant energy loss as well, and in fact (if the metal is burning in an infinite medium which acts as a black body receiver) this loss can be greater than the amount radiated to the metal surface. For metallized solid propeliants burning in an internal burning grain, there would be no radiant energy loss since the energy would be radiated back to the propeilant. For metals strands, however, burning in strand bombs, there would be a definite radiant energy loss.

Now that the distinction between metal diffusion flames and inquid droplet flames has been made, certain physical characteristics of metal combustion will be examined.

III. THE METAL FLAME TEMPERATURE

Brewer (2) reports that under highly oxidizing conditions to almost neutral conditions the metal oxide molecule (represented as M_XO_Y) is the main vaporizing species for all the alkaline earth oxides except BeO, which decom-

elements even under neutral conditions and under reducing conditions all the alkaline earth oxides vaporize predominately to the metal atom vapor. From the description of the diffusion flames given earlier, it was pointed out that for analysis purposes the fuel and oxidizer merge in some proportion and react in a definite front. More realistically, the two begin to react prior to reaching the analytically prescribed proportions and do react whenever a ratio of fuel vapor to oxidizer suitable for reaction has been reached. The largest majority of the metal then is reacting under reducing conditions. Since the metal oxide which is formed would decompose to the elements as suggested by Brewer (2), higher concentrations of oxygen would be found which further prevents the diffusion of oxygen to the reaction zone. Thus, it appears that the reaction of metal vapors in diffusion flames will be under reducing conditions and the product oxide formed will decompose upon vaporization. Decomposition upon vaporization then limits the flame temperature to the vaporization temperature of the oxide.

For boron, Brewer (2) shows that from the spectroscopic dissociation energy of BO that B2O3 cannot vaporize as BO. However under reducing conditions B_2O_3 will vaporize as BO. For example, a mixture of B and B_2O_3 should produce a Latmosphere partial pressure of BO at 1100° K. BO is stable enough that even the most stable oxides can be reduced by boron under vacuum conditions. However, the B - B_2O_3 condition is that which exists in elemental boron diffusion flames and here too evidence shows that the flame temperature must be the vaporization temperature of the oxide B_2O_3 .

Brewer and Searcy (la) investigated the Al - Al_20_3 system in an effusion cell. From the vapor pressure measurements they obtained a boiling point of $3800^{\circ} \pm 200^{\circ}$ K for Al_20_3 . Farber (26) reports that their evidence for this dissociation of Al_20_3 to Al_20 and AlO at high temperature has been substantiated by others. Al_20 is the product in a reducing atmosphere and AlO in an exidizing atmosphere. From thermodynamic calculations, Farber was able

siderable. Here too then is the limiting vaporization temperature effect. Farber also showed thermodynamically the extent of decomposition of Al_2O_3 under various atmospheres and conditions which can exist in propellant flames. In the presence of CO and H_2 , Al_2O_3 starts to dissociate appreciably at temperatures about 3000°K.

However, the author believes that these reactions with CO and H₂ are only of import in later stages of propellant flames when the aluminum particles have burned already as diffusion flames. The diffusion of CO and H₂ to the burning metal particle will not be as great, of course, as the diffusion of oxygen. Small equilibrium concentrations of those species will be built up through the complete diffusion flame zone. More likely the important reducing reaction would be

2 Al+ Al203 -2Al20 + 0

Farber lists the value of \triangle H_f° (25°C) for Al₂O(g) as -34 kcal/mole. Since the \triangle H_f° of exygen atoms is 59 kcal/mole the decomposition-vaporization reaction consumes about all the energy of formation of the Al₂O₃ and the temperature must be limited as much. Since BO has a \triangle H_f° of only -5.3 kcal/mole the same situation exists for boron. It is obvious for metal exides which decompose into elemental metal vapor and exygen atoms.

Thus, until evidence is put forth to the contrary, it appears safe to establish the vaporization temperature of the metal vapor oxide as the adiabatic flame temperature in a metal diffusion flame for large groups of metals. The groups which fall in this category are:

- (a) Alkali metals, which include Na, K, Rb, and Cs. Li which is also in this group may have a stable gaseous oxide. However, this is yet to be determined positively.
- (b) Alkaline earth metals, which include Be, Mg, Ca, Sn, and Ba.

- (c) Boron-aluminum group which includes B, Al, Sc, Y, and La.

 Gaseous sub-oxides exist for these groups.
- (d) Titanium zirconium
- (e) Manganese
- (f) Iron group, which includes Fe, Co, and Ni.
- (g) Zinc group, which includes Cd also.
- (h) Silicon

As more data appear, each metal-oxygen system must be evaluated independently.

Table II taken from Grosse and Conway (42) lists properties of the metals

and metal oxides and the melting and boiling points. Consequently, the boiling

point temperature of the oxide may be taken as the combustion temperature.

The dissociation reactions which limit temperatures obtainable are all affected by the pressure. Naturally, as the pressure is raised the dissociated products will reassociate as given by the Le Chatelier's principle. Examination of the equilibrium constants for the vaporization dissociation step readily shows that the vaporization temperature increases with an increase in pressure. It is estimated that at 1000 psia the temperature of burning aluminum would be about 5900°K.

IV. BURNING CHARACTERISTICS

With the establishment of the criterion that the boiling point of metal oxide is the metal flame temperature, certain other conclusions as to the character of the metal diffusion flame may be drawn. These conclusions are based on certain physical properties of the metals themselves.

If the vaporization temperature of the oxide is greater than the vaporization temperature of the pure metals, then the structure of the diffusion flame would be much like the double film concept applied to liquids. However, if the vaporization temperature of the pure metal is greater than

TABLE II

NORMAL MELTING, AND BOILING POINTS OF METALS AND THEIR HIGHER OXIDES*

Element	Element Melting Point, T _M O	Element Boiling Point, T _M OK	0×i de	∆H ^O oxide kcal/mole	Oxide Melting Point TMOK	Oxide Boiling Point T V OK
Li	460	1640	Li ₂ 0	-146.6	2000	2600
В	2570	2820	B ₂ 0 ₃	-305.4	720	2520
Na	370	1150	Na ₂ 0	- 9 9.4	1193	1550
M g	920	1370	MgO	-143.7	3075	3350
Al	930	2720	A1203	-400.2	2323	3800
κ	330	1030	κ ₂ 0	- 86.4	800	1750
Ca	1120	1960	Ca0	-151.7	2860	3800
Ti	2000	3530	Ti ₂ 0 ₃	- 363.0	2400	3300
Zr	2120	5770	Zr0 ₂	-261.8	2960	4570
Be	1550	3240	Be0	-143.7	2820	4530
Si	1700	2870	sio ₂	-210.2	2000	2500

^{*} Considering the basic uncertainties in the thermodynamic and physical data at high temperatures, the author wishes to point out that the data presented throughout this paper are the best he could find in today's literature. Since most data were used for relative comparisons, no effort was made to give statistical deviations from the values reported. Such deviations in many cases can be large.

the vaporization temperature of the oxide, then reaction cannot take place in the vapor phase but must take place on the moiten surface at some temperature lower than the vaporization temperature of the metal but about that of the vaporization temperature of the oxide. The approach used to calculate the case when $T_{MO} > T_{M}^{V}$ is somewhat similar to that used for liquids and the distinctions between the two were discussed in the previous section. From Table II. It is found that Li, Na, Mg, Ai, Ca, K, Be, and Si fail into this group.

When $T_M^V > T_{MO}^V$ then the oxide accumulates on the molten metal surface. The procedure for calculating the burning rate in this case becomes much more complex. The rate may be determined by the surface chemistry or by the diffusion rate of the oxidizer through a solid oxide film. B, Ti, and Zr fall into this group. The character of the film can be different according to the physical properties of the oxide. If the oxide has a lower density than the moiten metal, then it will float on the surface and oxygen must reach the raw metal by crevices or openings formed due to oxide running off the surface, cracking due to expansion of the liquid, etc. However, if the oxide is lighter than the molten metal, but soluble in the metal, then some metal would always appear on the surface. If the oxide is more dense than moiten metal, it would sink through the moiten surface and keep a metal surface exposed to the oxidizer. Physical Intuition would suggest perhaps that the burning rates of the metals would fall generally in categories according to the descriptions given above; i.e., those metals whose oxides are heavier than the molten meta! would be the fastest; those which are lighter and soluble, next; and those which are lighter but not soluble, the slowest. Further metals which fall in the case $T_M^V > T_M^V$ should burn faster than those in which $\mathsf{T}_\mathsf{M}^\mathsf{V}\!\!>\!\!\mathsf{T}_\mathsf{M}\!\!\circ\!\!\mathsf{N}$. Thus, under the same conditions Al and Mg should exhibit faster burning rates than B, Ti, or Zr. Because the physical picture of metals in which $T_{MO} > T_{M}$ is the only clear one, analytical determination of the burning rates of this group alone will be treated here.

The first analytical attempt to determine the burning rate of metals taking into consideration that the oxide vaporization temperature establishes the flame front temperature was due to Coffin (77). Coffin, however, neglected the possibility of radiant heat transfer processes. The calculation of burning rates, when radiation is included, is extremely difficult. Nevertheless, the form of the resulting expressions can indicate the importance of various physical and chemical parameters. Given below is the development of Coffin and Brokaw (37) for the burning of magnesium; this development is modified to take into account radiation effects.

Coffin replaces the flame front of Infinitesimal thickness used in the regular double film concept by a high temperature zone of diffusion and reaction at the boiling point of the oxide. Pictorially this model is reported in Figure 4 and could be thought of as three films. MgO condenses outside the high temperature zone and it is assumed that this product cannot diffuse toward the droplet.

The high temperature zone BB' is one of constant temperature. The oxygen concentration is zero at B and some unknown value of B'. Inside the zone MgO can exist and dissociate to the elements. At B', MgO condenses completely.

In the AB region, there is not net flow of any component except fuel. Products are assumed to pass outward. The energy equation in zone AB is then

where W_f = fuel mass transfer rate, moles/sec

△H = sensible heat of fuel to the bolling point, cal/mole

LV = latent heat of vaporization, cal/mole

H = enthalpy, including chemical energy, cal/mole

X = radius from center of particle, cm.

k = thermal conductivity, cal/cm sec °C

T = temperature, °C or °K

RG = a radiant energy term

The only difference between this equation and that of Coffin and Brokaw is the term Rg. Rg is a function of the emissivities $(\mathcal{E}_A,\mathcal{E}_B)$ and the absorbing characteristics of the Mg vapor in AB. For a droplet system the area S at A and at B enter in the radiation term, however, in a strand burning along the cylindrical axis Equation I would be developed by a one-dimensional approach and the area cancels throughout all equations.

For a droplet, though, and particularly for most metals, S_B - S_A is so small that S_B = S_A and thus the solution of the final equations is simplified. The actual conditions in BB' are such that E_B =I, as discussed in Section II. E_A is affected by the reflectivity of the surface, and if the molten surface is highly reflective E_A could be of the order of a tenth. Nevertheless T_B 4 - T_A 4 is so large, at high values of pressure and oxygen particle pressure that it would appear the radiation term still could not be neglected. Thus, the physical condition of the surface of molten metal can be of some importance in determining its burning rate.

In the region B', oxygen diffuses inward, and the MgO floats out. The energy equation in this zone is

$$\omega_{\xi} \left[Q + m_{0} \left(H_{0} - H_{0} \right) - m_{M_{0}} \left(H_{M_{0}} - H_{M_{0}} \right) \right] = 4\pi \kappa^{2} k_{B_{c}} \frac{d7}{dx} + R_{L}$$
 (2)

where Q = heat of combustion at reference temperature, cal/mole n = stoichiometric coefficient

RL = a radiant energy term

R_L is a term which enters only in droplet or strand burning tests and is a radiation loss to the surrounding.*

^{*} The introduction of the terms R_G and R_L is similar to the suggestion of Summerfield (g) that such terms should be introduced in determining the burning rates of composite propeliants.

Coffin and Brokaw do not include this RL term either. If the particles are burning in an internal burning solid propellant grain, the energy is transmitted back to the grain and is not lost. Since most basic experiments are performed with spheres or strands, this term should be included when comparisons are made between experiment and theory. Since the energy is accepted by an infinite absorbing atmosphere or cold walls, R₁ can be written, then, as

$$R_{L} = \sigma S_{A} T_{B}^{Y}$$
 (3)

where σ is the Stefan-Boltzmann constant. Again for this term it is assumed that $S_{B^{\dagger}} \cong S_{B} \cong S_{A}$.

Because of the condensed product, the diffusion equation for oxygen generally used for other systems reduces to the equation for the diffusion of oxygen through a stagnant film of ambient inert gas

$$\frac{W_{o_2}}{4\pi\chi^2} = -\frac{D_{o_2}P}{RT(P-p_{o_2})} \frac{dp_{o_2}}{d\chi}$$
(4)

where D = diffusion coefficient, sq cm/sec

P = total pressure, atm

p = partial pressure, atm

R = universal gas constant

The diffusion equation for oxygen and the energy equation for the outer zone B¹C, combined and integrated, is

$$\frac{D_{o_{3}}}{m_{a}RTk_{s'c}}\int_{P_{o_{3}}}^{P_{o_{3}}}\frac{P}{P-p_{o_{3}}}dP_{o_{3}}=\int_{T_{o_{3}}}^{T_{o_{3}}-T_{o_{3}}P}\frac{dT}{(Q+m_{o_{3}}(H_{o_{3}}-H_{o_{3}}c)-m_{m_{o_{3}}}(H_{m_{o_{3}}}-H_{m_{o_{3}}}c)-m_{m_{o_{3}}}H_{m_{o_{3}}}c}}{(D+m_{o_{3}}(H_{o_{3}}-H_{o_{3}}c)-m_{m_{o_{3}}}(H_{o_{3}}-H_{o_{3}}c)-m_{m_{o_{3}}}H_{m_{o_{3}}}c}-\frac{R_{o_{3}}}{W_{o_{3}}}$$
In the high temperature zone BB' the energy equation as written

by Coffin Is

$$W_{m_g} H_{m_g} + W_{o_2} H_{o_2} + W_{m_g} O H_{m_g} O - W_{5} H_{5} = 0$$
 (6)

and the continuity for each species is

$$2 W_{02} + W_{m_00} = 0$$

$$W_{mg} + W_{mg0} = W_{f}$$
 (8)

Combining with the energy equation

$$W_{m_{g}}(H_{50}-H_{m_{g}})-2W_{02}(H_{m_{g}0}-H_{50}-\frac{1}{2}H_{02})=0$$

$$W_{m_{g}}=-2W_{02}\left(\frac{H_{m_{g}0}-H_{40}-\frac{1}{2}H_{02}}{\Delta H+L^{2}+H_{m_{g}}-H_{m_{g}}A}\right)$$

$$W_{m_{g}}=-2\theta W_{02}$$
(9)

 $H_{MgO} - H_{fo} - \frac{1}{2} H_{O_2}$ equals the heat of combustion of Mg at T_O to form condensed MgO at the boiling point for the given conditions. Then

$$\mathcal{W}_{\varphi} = -2(\Theta + I) \, \mathcal{W}_{o_2} \tag{10}$$

At atmospheric pressure Θ for magnesium equals about two. It may be regarded as the ratio of the amounts of MgO condensed at the outer and inner boundaries of the BB' zone.

With substitution of the above equation into the diffusion equations one has

$$\frac{\omega_{\pm}}{4\pi\chi^{2}} = \frac{D_{0}P(2\theta+1)}{RT[P+(2\theta-1)p]} \frac{dp_{0}}{dx}$$
(11)

Integrating the equations one obtains

$$\int_{\chi_{A}}^{\chi_{B}} \frac{W_{+}}{4\pi\chi^{2}} d\chi = \int_{T_{A}}^{T_{B}/} \frac{R_{AB}}{[H_{n_{g}} - H_{f_{g}} - \frac{R_{G}}{N_{f}}]} dT = \infty$$

$$\int_{\chi_{B}}^{\chi_{B}} \frac{W_{+}}{4\pi\chi^{2}} d\chi = \int_{\rho_{a}=0}^{R_{a}/B} \frac{D_{o_{2}}P}{RT} \frac{(2\theta+2)}{\rho_{+}(\partial\theta-1)\rho_{*}} d\rho_{a} = \beta$$
(12)

Combining the above three integrals

$$\int_{k_A}^{\gamma_c} \frac{W_f}{4\pi \gamma^2} d\chi = \alpha + \beta + \delta$$
 (15)

For $\frac{1}{\lambda} \rightarrow 0$

$$\frac{W_{5}}{4\pi T X_{A}} = \alpha + \beta + \delta \tag{16}$$

A possible approach to the solution of the equations would be to assume a value of W_f so that the radiation terms could be calculated and then determine the W_f from the remainder of the equations. Reiterate until the value assumed is equal to the value calculated.

 \mathcal{L} , \mathcal{B} , and \mathcal{F} are positive with the log term dominating, particularly in oxygen rich atmospheres. Equation 5 used for calculating P_{2} , \mathcal{B}' shows a strong dependence of P_{2} , \mathcal{B}' upon P_{2} , \mathcal{C} . Further, it has been shown that the vaporization temperature of the metal oxide is dependent on partial pressure of the oxygen present, that is P_{2} , P_{2} . This dependence is felt strongly in the radiation term introduced where the temperature is raised to the fourth power. Thus, for metal combustion there is an even stronger dependence of burning rate of the oxygen partial pressure than there is in liquid droplet combustion.

In the above analysis it was assumed that none of the product oxide diffused toward the droplet. There is some question as to how good this assumption is for various substances. Obviously, the droplet represents a

there could be deposition of oxide on the moiten metal surface, even in the class of metals discussed here, and thus the burning rate can be hindered in the manner described for the metals which surface react. Another factor which could determine the amount of oxide that would diffuse back to the moiten surface and thus hinder the burning rate, is the relative location of the burning front to the metal surface. Those metals which would have fronts closer to the surface than others would be expected to have the greatest difficulty from oxide deposition.

From the above analysis It would follow then that aluminum and iithium would show more serious oxide surface deposition due to diffusion than magnesium. One could estimate this relatively from the ΔH_{TeT} data in Table III and through discussion given above.

TABLE III
HEAT CONTENT TERMS FOR VARIOUS MATALS

Element	Cp Solid	M.P T ^o C	∆ H Solid Cal/gm	L M Cal∕gm	C _p Ii q uid	BP T ^O C	∆ H liquid cal/gm	L ^V cal/gm	∆H _{ror} cal/gr	T _{MO} °C
Ве	.50					3170	1570*	8355 *	9930	4000
Li	1.10	186	177	32.8	1.1	1317	1245	4648	5100	2300
AI	.24	660	152	95.4	.26	2500	464	2260	2970	3500
Mg	.28	651	I 75	88.8	.33	1103	150	1337	1750	3100
Na	.28	98	20	27.2	.30	883		1005	1290	1300
с ₆ н ₆					.40	80	22	94	116	

^{*} Based on sublimation process

V. IGNITION

The balance between a rate of heat production and the rate at which heat can be removed determines whether a substance will ignite and react (46). In the case of fuel droplets and particularly for metals, this factor is largely dependent upon how fine is the sub-division, i.e. particle size. Smaller particles have a large surface to volume ratio, and hence a higher rate of heat production per unit volume, and a lower cooling rate.

If the surface area to volume ratio is very large, the conduction of heat from the flame front by the metal particle is greatly reduced. Heat loss then is by the relatively slow heat transfer process in a gas. Thus, neighboring particles of metal reach the ignition temperature with great rapidity. If finely divided powder is loosely packed together, the heat loss by convection largely is eliminated and under such conditions of low heat loss, spontaneous ignition of metal can occur.

When solid metal rods are used, conduction by the metal rapidly absorbs large quantities of heat from the burning front, and little heat is left for raising the temperature of the adjacent metal to the ignition temperature.

Consequently, as with other combustible substances, when in bulk, metal can be made to burn at a controlled rate, whereas when divided, it will burn with great violence.

In regard to metal particles of widely different sizes flowing in a combustible stream, it would appear then that large particles will take a longer time to heat to the ignition point, and thus have longer ignition time than smaller particles. Thus, there will be a definite size effect in regard to ignition.

The character of the Ignition process is entirely different from that of the combustion process in that the ignition process is most likely not a diffusion type flame, but a surface burning phenomena. The transition

pattern from a surface ignition process to a diffusion flame has not been reported. The process is most likely something as follows. The surface reacts and the temperature rises sufficiently to vaporize the metal and subsequent reaction between the hot metal vapor and oxygen proceeds as a diffusion flame. However, if during the ignition process the oxide coats the surface so that there is slow burning of the ignited metal, then radiative transfer could vaporize and ignite neighboring particles which would burn as diffusion flames. However, if the ignition sources were energetic enough to vaporize the metal initially, there is no doubt the metal would burn as a diffusion flame.

Grosse (42) reports that metals can be divided into three classes according to their ignition process. In the first class are those metals that ignite below their melting point and includes Mg, Ca, Ti, Zr, Mo, Fe, Ca; in the second class are those that ignite above their melting point and include Al, Li, and Na; and, of course, there are those metals that do not ignite, Hg, Ag, and Pt.

VI. DISCUSSION AND CONCLUSIONS

Throughout the previous sections reference was made to various practical conclusions that would be reached by the description of the metal burning process proposed here. A more detailed discussion of the various topics will be given here.

One topic of major interest worth further discussion is the temperature reached in the burning of metals. As previously stated most metals of interest in propulsion are limited to combustion temperatures which correspond to the vaporization temperature of the higher metal oxide. However, this temperature varies with the pressure or in the case of a mixture of gases with the oxygen, or oxidizer, partial pressure. In propulsion systems

when the metal is only a fuel additive, a somewhat more complex situation exists. Here there are two competing processes. One is associated with the burning of the fuel and the other with the burning of the metal. It is proposed that the metal particles burn as diffusion flames described in earlier sections of the article. If this is so, the temperature of the diffusion flame zone surrounding the burning metal particle will be dependent upon the oxygen partial pressure and this temperature will be different from the adiabatic combustion temperature of the complete propellant mixture. Except under unusual circumstances one would expect the diffusion flame temperature of the metal to be greater than the adiabatic flame temperature of the mixture. The equilibration of temperature zones must take place outside the active reaction zone. The availability of oxygen for the metal combustion in a solid propeliant is difficult to predict. If the oxidizer and hydrocarbon fuel are also burning as a diffusion flame as is most likely in a composite propellant (q), then the metal should not find too great a deficlency of oxygen. Even in homogeneous propellants the various reaction states taking place would be slow enough that any metal particles added would be able to compete for oxidizer. However, one could think of many situations, such as fine metal particles added to a premixed hydrogen-oxygen gas mixture prior to entering a flame front, in which the gaseous fuel would burn with much greater rapidity than metal. The metal would thus burn in a stream depleted of oxygen and the temperature of the flame around the metal particle would be lower than if It had burned more completely In the earlier stages of the flame.

Obviously then, when the metal is an additive to a <u>fuel</u> or a solid propellant the situation is most complex. However, in composite solid propellant it would appear to the author that the temperature of the flame surrounding the metal particle would be much greater than the equilibrium temperature and that the possibility that radiant energy exchange is an important factor

in determining the burning rate, as suggested by Summerfield (g), is an excellent one. Of course, the higher the temperature the more pertinent the radiation term. The higher the radiation term the less sensitive should be the pressure dependency of the burning rate.

The character of the burning process of substances which contain metallic elements will vary according to the reaction procedure for complete oxidization. Consider for an example a metal which falls in the category $T_M^V > T_{MO}^V$. Such metals when added in elemental form or through a compound which decomposes to give the elemental form of the metal should have long burning times and thus could give poor combustor efficiencies. Boron is a typical example. However, if the metallic element could be added in compound form such that the compound does not decompose before oxidization, then oxidization of the compound directly could take place and the elemental burning process eliminated. Again taking boron as an example, if a boron compound could be oxidized in the same manner as an ordinary hydrocarbon where the oxygen attaches to the carbon (or boron for the example) before decomposition (h) the slower burning rate of metal burning for $T_M^V > T_{MO}^V$ is circumvented.

For the alkylated compounds which do decompose into hydrogen, hydrocarbon fragments, and the elements, then there can be stages of combustion with the hydrogen burning first, fragments next, and then the metals. As discussed above, the depletion of oxygen concentration could effect the temperature of the metal flame front and the burning rate.

 less oxide deposition may be a great oversimplification. This area may be a fruitful one for experimental investigation. The author and his students are initiating such experimental investigations which will be based on burning of small diameter metal strands at high oxygen concentrations and total pressures. Pressure release during strand burning should quench the fixme and permit surface examination by metallurgical photomicrography. Fiame front and metal surface separation distance measurements will be attempted by optical techniques.

Other experimental investigations will include the measurement of burning rates of metal strands under various conditions in order to compare the results with the theory which was given earlier and which include the radiation terms. Relating strand rates to spherical particle rates has recently been accomplished by Mead (j). Spectroscopic examination of the various zones of the diffusion flame is anticipated and should give most important information on the metal combustion process.

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VIII. NOMENCLATURE

 $c_p = specific heat, cal/gm {}^{O}C$

 $D = diffusion coefficient, cm^2/sec$

H = enthalpy, including chemical energy, cai/mole

AH = sensible heat of fuel, with no subscript then to the boiling point, cal/mole ΔH_C = heat of combustion, Btu/lb fuel or Btu/lb (fuel and oxldizer)

ΔHm= ΔHsolid + LM + ΔHIIquid + LV, cal/gm

k = thermal conductivity, cal/gm cm °C

Kp = equilibrium constant

L = latent heat, cal/mole or cal/gm

n = stoichiometric coefficient

In = total number of moles of gases

p = partial pressure, atm

P = total pressure, atm

Q = heat of combustion at reference temperature, cal/mole

R = universal gas constant, cm³ atm/ mole

 R_{G} , R_{L} = radiant energy terms

S = surface areas, cm²

T = temperature, °C or °K

W = mass transfer rate, moles/sec

X = radius from center of particle, cm

a, B, 8 = integrals

E = emissivity

 $\sigma = \text{Stefan-Boltzman constant}$

 θ = heat content ratio

Subscripts

A, B, B', C = positions specified in Figures 3 and 4

f = fuel

I = Ilquid

M = metal

MO = metal oxide

o = reference state

Superscripts

M = melting point

V = vaporization or boiling point

3

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XI. LIST OF FIGURES

- Figure 1 Heat of Combustion in Btu per pound of fuel and oxidizer vs atomic number of principal elements, from Carpenter (C).
- Figure 2 Heat of combustion in Btu per pound of fuel vs atomic number of principal elements, from Carpenter (C).
- Figure 3 Cross-sectional model of burning fuel particle, from Coffin and Brokaw (37).
- Figure 4 Cross-sectional model of burning fuel particle having high boiling point exhaust product, from Coffin and Brokaw (37).

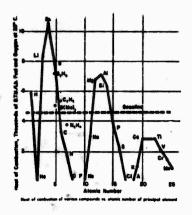


FIGURE I

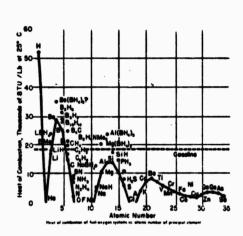


FIGURE 2

